

From QFT on noncommutative spaces studied with Manfred to integrable QFT models

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Personal prehistory

In January 2000 I joined [Harald Grosse](#) at Universität Wien. This was my 2nd postdoc after CPT Marseille.

My plan was to resolve a puzzle discovered by [Minwalla, van Raamsdonk, Seiberg, 1999]:

Ultraviolet/Infrared-mixing

Certain conditionally UV-convergent (but absolutely divergent) subgraphs in QFT-models on [noncommutative \$\mathbb{R}^D\$](#) produce [untreatable IR-divergences](#).

I had some hope that gauge models are better. With Harald and Thomas Krajewski we quickly proved that the [ghost sector is fine](#), but in the end this [did not help](#).

Next hope was [SUSY](#), where quadratic divergences cancel. We did not have much experience with SUSY.

First contact with Manfred

It was great luck that **Manfred** contacted Harald precisely during this time of uncertainty.

Manfred, with a group of young students, was already working on **SUSY on noncommutative \mathbb{R}^D** .

He invited us to join – the beginning of a very successful collaboration.

In fact the results were ready, but Harald and I wrote the introduction with an extended overview of the field:

“The superfield formalism applied to the noncommutative Wess-Zumino model,” hep-th/0007050.

Manfred's work with students

Manfred always had very many students.

He had a remarkable method to guide them:

Manfred regularly scanned arXiv. When he found an interesting hot topic, he gave it to a student, or a group of students.

He had an exceptional intuition what topic is fruitful.

Students

I had interaction with 3 generations of students:

- 1 The closest with [Andreas Bichl](#), [Jesper Grimstrup](#) and [Lukas Popp](#) on the Seiberg-Witten map (more below).
- 2 [Volkmar Putz](#) joined me in Leipzig; we applied the Seiberg-Witten map to super-Yang-Mills.
Further coauthors: A. Gerholt, M. Ertl, M. Wickenhauser
- 3 Another generation around [Peter Fischer](#) worked on deformed Minkowski space.
Further coauthors: H. Bozkaya, M. Pitschmann.

Later generations include [Michael Wohlgenannt](#) and [Daniel Blaschke](#), and probably others in the audience

Moyal space = noncommutative \mathbb{R}^D

vector space of Schwartz functions, equipped with non-local, associative, noncommutative product

$$(f \star g)(x) = \int_{\mathbb{R}^2 \times \mathbb{R}^2} \frac{d\eta dk}{(2\pi)^2} f(x + \frac{1}{2}\Theta k) g(x + \eta) e^{i\langle k, \eta \rangle} \quad \Theta = \begin{pmatrix} 0 & \theta \\ -\theta & 0 \end{pmatrix}$$

- strict deformation by \mathbb{R}^D -group action (Rieffel)
- compared with formal deformation quantisation (Kontsevich)
- restricting formal deformation to finite order gives local theory

For gauge theory, [Seiberg-Witten, 1999] found a **map to higher-derivative gauge theory** on usual \mathbb{R}^D .

I was sceptical that this is interesting. The young students, after discussion with the group of Julius Wess, convinced me that we should try. It turned out to be a gold mine.

Seiberg-Witten map

Thanks to Mathematica programming skills available at TU Wien we were several months ahead of all other groups.

Our strongest result is:

Theorem [Bichl-Grimstrup-Grosse-Popp-Schweda-W, 2001]

After field redefinition, the Seiberg-Witten mapped photon selfenergy is renormalisable to all orders in θ .

In the very end this did not help: as soon as there are fermions, the method breaks down.

I abandoned the subject in summer 2002, also paused the interaction with Manfred

Final attempt: matrix basis and RG

UV/IR-mixing resisted all treatments. Now being at MPI Leipzig, I started with Harald a final attack based on:

- Use **matrix basis of Moyal space**. No positions, no oscillating phase factors, direct control of convergence (inspired by José Gracia-Bondía)
- Control large- N matrix limit by **renormalisation group** à la Polchinski (inspired by Thomas Krajewski)

Matrix basis (in 2D)

$$\phi(x) = \sum_{m,n=0}^{\infty} \Phi_{mn} f_{mn}(x)$$

$$f_{mn}(x) = 2(-1)^m \sqrt{\frac{m!}{n!}} \left(\sqrt{\frac{2}{\theta}} x_1 + i x_2 \right)^{n-m} L_m^{n-m} \left(\frac{2|x|^2}{\theta} \right) e^{-\frac{|x|^2}{\theta}}$$

satisfies $f_{mn} \star f_{kl} = \delta_{nk} f_{ml}$ and $\int dx f_{mn}(x) = 2\pi\theta \delta_{mn}$

This alone did not help, but it pointed out a **natural modification of the noncommutative $\phi^{\star 4}$ -action**

Harmonic propagation

Adding a 4th marginal coupling makes everything well-defined:

$$S[\phi] = \int d^4x \left(\frac{Z}{2} \phi \star (-\Delta + \Omega^2 (2\Theta^{-1}x)^2 + \mu_0^2) \phi + \frac{\lambda_0 Z^2}{4} \phi \star \phi \star \phi \star \phi \right)(x)$$

- **renormalisable as formal power series** in λ_0 [Grosse-W., 2004]
(renormalisation of μ_0^2 , λ_0 , $Z \in \mathbb{R}_+$ and $\Omega \in [0, 1]$)
means: well-defined **perturbative** quantum field theory
- Langmann-Szabo duality: theories at Ω and $\Omega^* = \frac{1}{\Omega}$ are the same; self-dual case $\Omega = 1$ is **matrix model**
- **one-loop RG-flow of running coupling is bounded**, no Landau ghost [Grosse-W., 2004]
- later: **β -function vanishes to all orders** in λ_0 for $\Omega = 1$ [Disertori-Gurau-Magnen-Rivasseau, 2006]

Manfred's help with my habilitation

- I tried to assemble renormalisation proof and one-loop vanishing of β_λ into a habilitation thesis.
- University of Leipzig rejected my request (but permitted lectures).
- Manfred rescued me and offered to do the habilitation at TU Wien (submitted 08/2004, defended 05/2005).
- Manfred also supported applications 2004+2005 for a Marie Curie Chair at TU Wien (was not realised; I got a job in Münster).

Main results since then (with Harald)

$\beta = 0$ should permit rigorous construction of $\lambda\phi^{*4}$

- action in matrix basis $S(\Phi) = V \operatorname{tr}(ZE\Phi^2 + \frac{Z^2\lambda}{4}\Phi^4)$
for $E = ((\frac{\mu_{\text{bare}}^2}{2} + \frac{4|\underline{n}|}{\theta})\delta_{\underline{n}\underline{m}})$, $\underline{n} = (n_1, n_2)$, $V = (\frac{\theta}{4})^2$
- use partition function $\mathcal{Z}(J) := \frac{1}{\mathcal{Z}(0)} \int D\Phi \exp(-S(\Phi) + \sum \Phi_{mn}J_{mn})$ as **tool to derive equations of motion (Schwinger-Dyson eq.)**
- covariance under $\Phi \mapsto U^*\Phi U$ yields Ward-Takahashi identity:

Theorem [Disertori-Gurau-Magnen-Rivasseau, 2006]

$$0 = \sum_{n \in I} \left(\frac{(E_a - E_p)}{V} \frac{\partial^2 \mathcal{Z}}{\partial J_{an} \partial J_{np}} + J_{pn} \frac{\partial \mathcal{Z}}{\partial J_{an}} - J_{na} \frac{\partial \mathcal{Z}}{\partial J_{np}} \right)$$

formula for $\sum_{n \in I} \frac{\partial^2 \mathcal{Z}}{\partial J_{an} \partial J_{np}}$ after controlling the kernel

Schwinger-Dyson equations

In a scaling limit $V \rightarrow \infty$ and $\frac{1}{V} \sum_{p \in I}$ finite, we have:

1. A closed non-linear equation for $G_{|ab|}$ (2009/2011)

$$G_{|ab|} = \frac{1}{E_a + E_b} - \frac{\lambda}{(E_a + E_b)} \frac{1}{V} \sum_{p \in I} \left(G_{|ab|} G_{|ap|} - \frac{G_{|pb|} - G_{|ab|}}{E_p - E_a} \right)$$

2. For $N \geq 4$ a universal algebraic recursion formula (2012)

$$G_{|b_0 b_1 \dots b_{N-1}|} = (-\lambda) \sum_{l=1}^{\frac{N-2}{2}} \frac{G_{|b_0 b_1 \dots b_{2l-1}|} G_{|b_{2l} b_{2l+1} \dots b_{N-1}|} - G_{|b_{2l} b_1 \dots b_{2l-1}|} G_{|b_0 b_{2l+1} \dots b_{N-1}|}}{(E_{b_0} - E_{b_{2l}})(E_{b_1} - E_{b_{N-1}})}$$

- similar formulae for other topological sectors
- no index summation in $G_{|abcd|} \Rightarrow \beta\text{-function zero!}$

Solution of the closed equation for $G_{|ab|}$

All **QFT-difficulties** (divergences, non-linearity, summation of perturbation series) **contained in a single equation for $G_{|ab|}$**

- becomes **integral equation** in limit $V \rightarrow \infty$ and $\frac{1}{V} \sum_{p \in I}$ finite
- **renormalisation of all Feynman graphs at once achieved**
- techniques from complex analysis permit reduction to fixed point equation for $G(x, 0)$:

$$G(x, 0) = \frac{1}{1+x} \exp \left(-\lambda \int_0^x dt \int_0^\infty \frac{dp}{(\lambda \pi p)^2 + \left(t + \frac{1 + \lambda \pi p \mathcal{H}_p[G(\bullet, 0)]}{G(p, 0)} \right)^2} \right)$$

- proved existence of solution with $G(x, 0) \xrightarrow{x \rightarrow \infty} \frac{1}{(1+x)^{1-\eta}}$,
conjecture $\eta = -\frac{\arcsin(\lambda \pi)}{\pi}$
- reflection positivity requires $\eta > 0$, i.e. negative λ

State of the art

- [Grosse-Sako-W, 2016] derives explicit and **exact formulae**, analytic in λ , **for all matrix correlation functions of $\lambda\phi^3$** .
- Perfect match with perturbation series
- $\lambda\phi^4$ currently on hold because exact solution of fixed point equation not yet available (but a lot is known).
- The topic inspired Razvan Gurau and Vincent Rivasseau to develop **coloured tensor models** as one approach to quantum gravity.
- We currently work on **translating matricial correlation functions back to Schwinger functions in position space**; main question concerns Osterwalder-Schrader axioms.

Exchange with Manfred was indispensable foundation for these achievements.